

# Quantum correlations with indefinite causal order

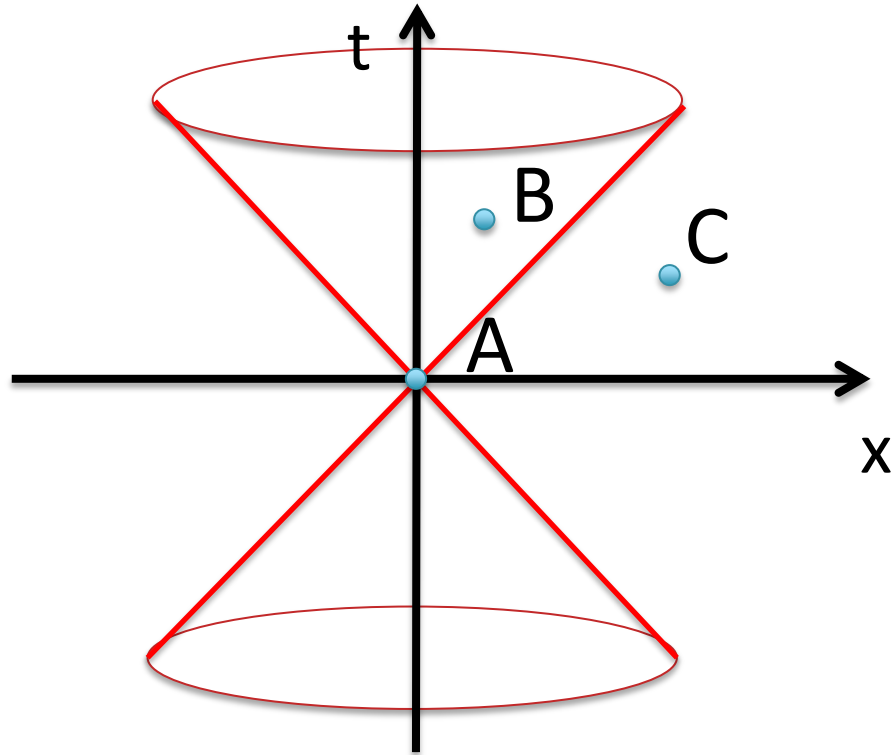
Ognyan Oreshkov, Fabio Costa, Časlav Brukner

Nature Communications (2012), doi: 10.1038/ncomms2076

(See also News & Views: Nature Physics 8, 860–861 (2012))

Emergent Quantum Mechanics, Vienna, October 5<sup>th</sup>

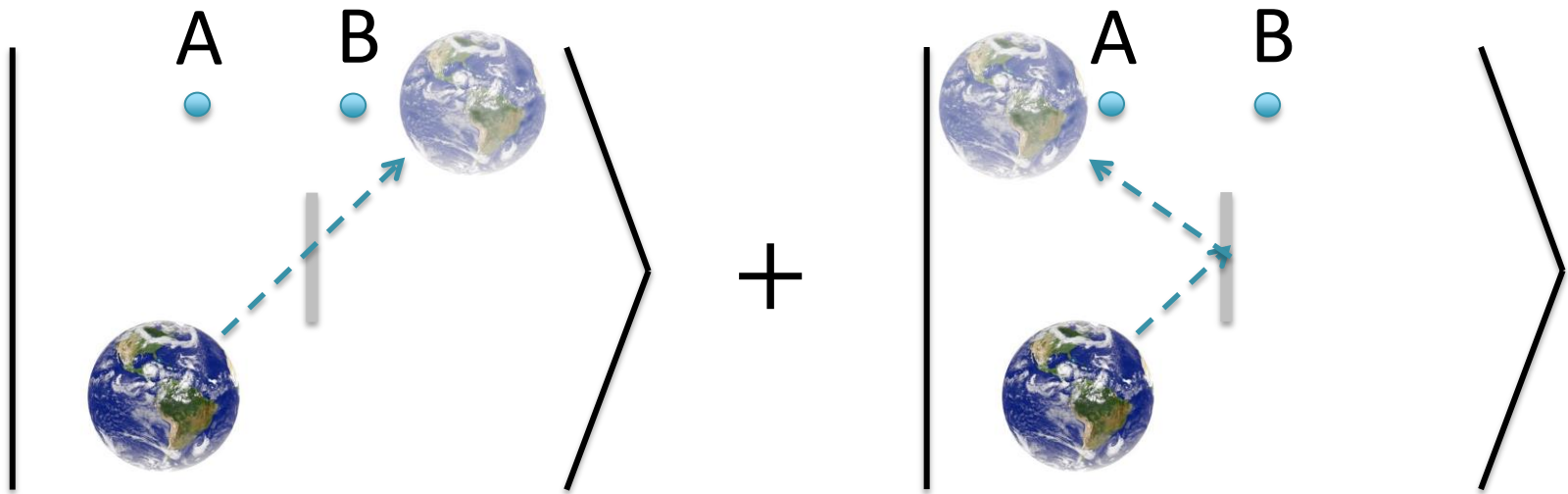
# Causality & space time structure



$A < B$ : A before B

$A \sim C$ : A and C causally neutral

# Quantum causal relations?



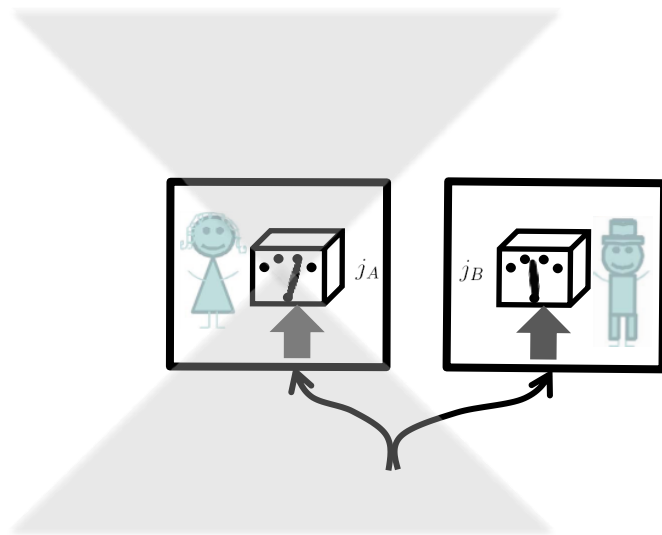
Space-time distance between A and B is not well-defined.

# Outline

- Framework for quantum mechanics with no assumed global causal structure:
  - includes all causally ordered (spatial and temporal) situations: shared states, channels, channels with memory, and probabilistic mixtures of these.
- Correlations that defy causal order:
  - Violation of a “causal inequality” – a communication task that cannot be accomplished with causally ordered operations.
- In a classical limit, causal order always arises → space-time may emerge from a deeper structure.

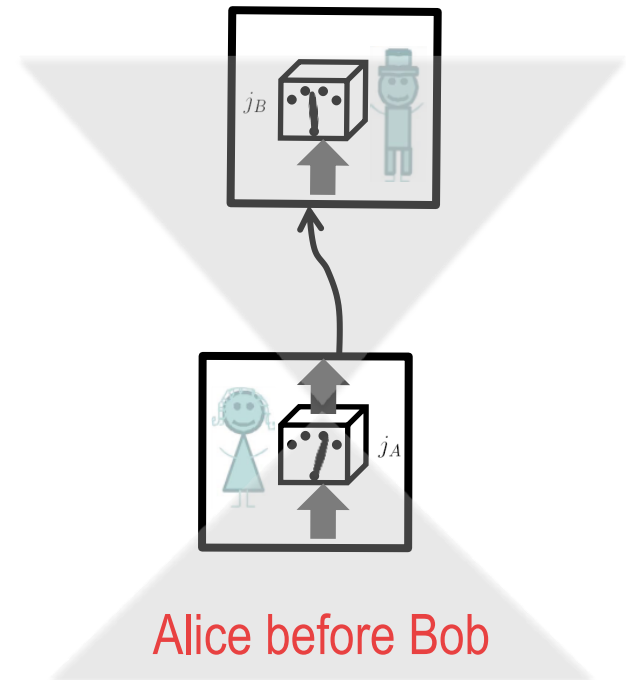
# Causal order from space-time

No-signalling



Space-like separated  
Causally neutral

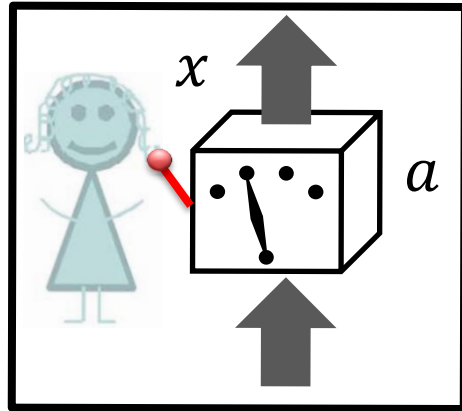
One-directional  
signalling



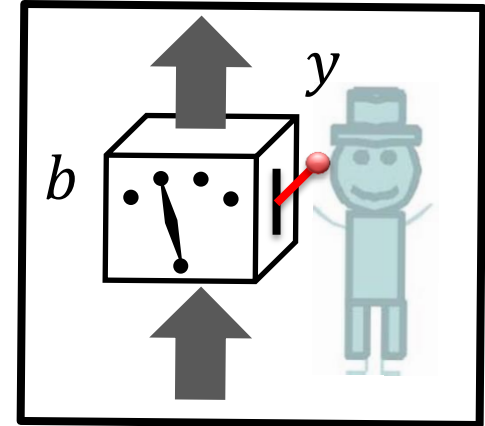
Alice before Bob

Time-like separated

# Operational meaning of signalling

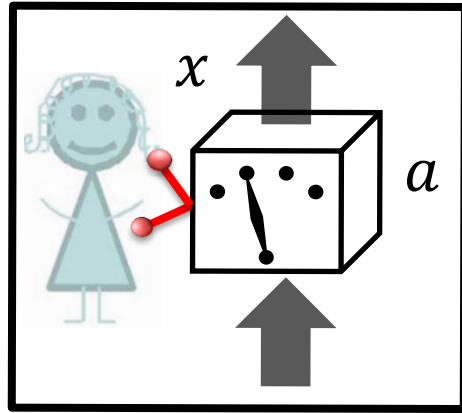


$$p(a, b | x, y)$$

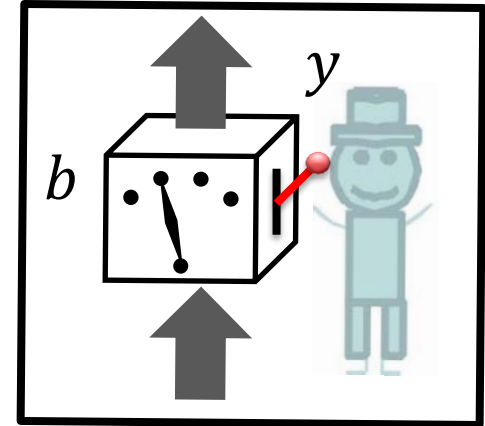


- $x$  and  $y$  are measurement settings (**statistically independent variables**)
- $a$  and  $b$  are measurement outcomes

# Operational meaning of signalling

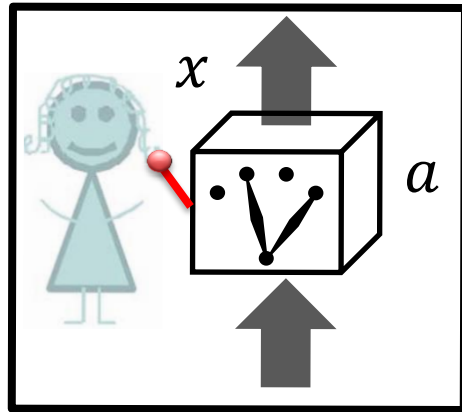


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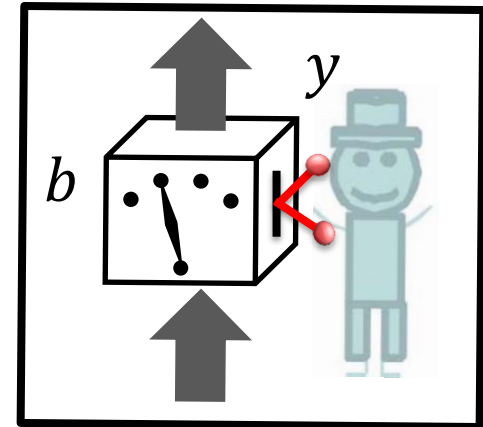


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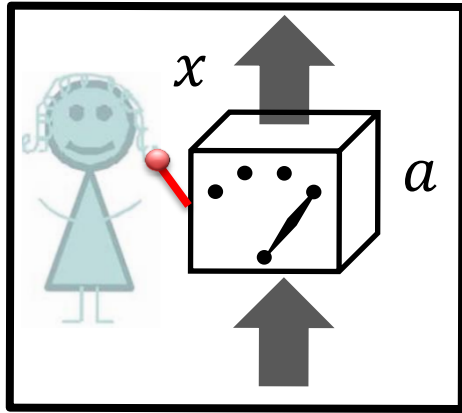
$$p(a, b | x, y)$$



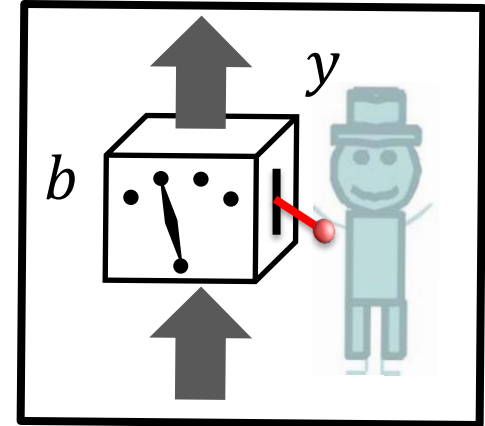
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# Operational meaning of signalling



$$p(a, b | x, y)$$



**No-signalling**

**One-directional signalling**

$$\sum_a p(a, b | x, y) = p(b | y)$$

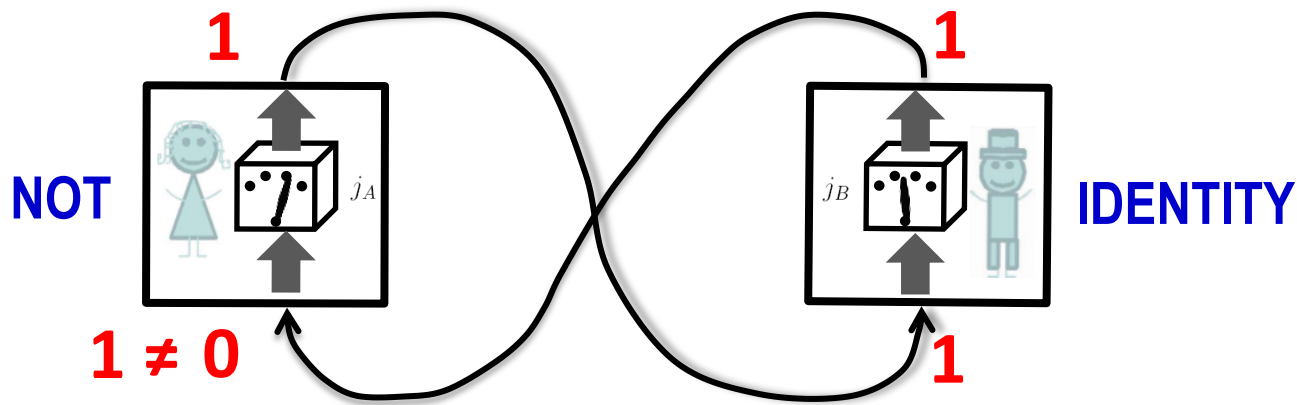
$$\sum_a p(a, b | x, y) = p(b | y)$$

$$\sum_b p(a, b | x, y) = p(a | x)$$

$$\sum_b p(a, b | x, y) = p(a | x, y)$$

# Two-directional signalling?

## Causal Loop



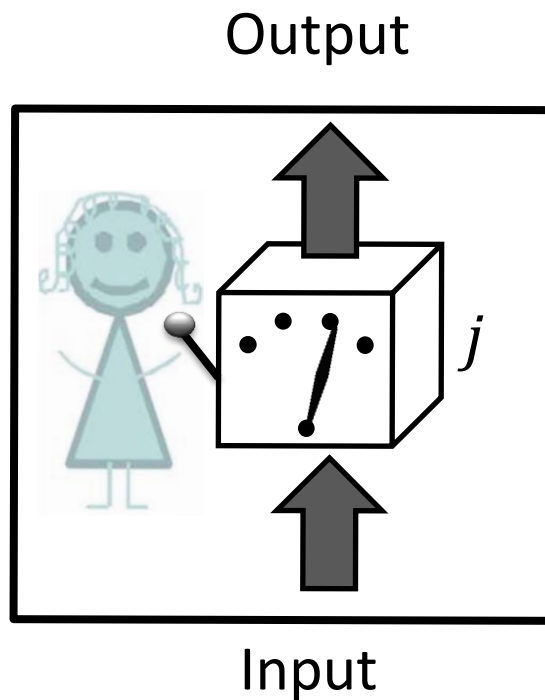
**Grandfather paradox!**

Gödel Universe: **Closed time-like curves** (CTC)

Proposed quantum solutions: Deutsch's or the Bennett-Schumacher-Svetlichny-Lloyd CTC-like structures are **non-linear** extensions of quantum theory

Linear structures, free of paradoxes?

# The framework: Closed laboratory



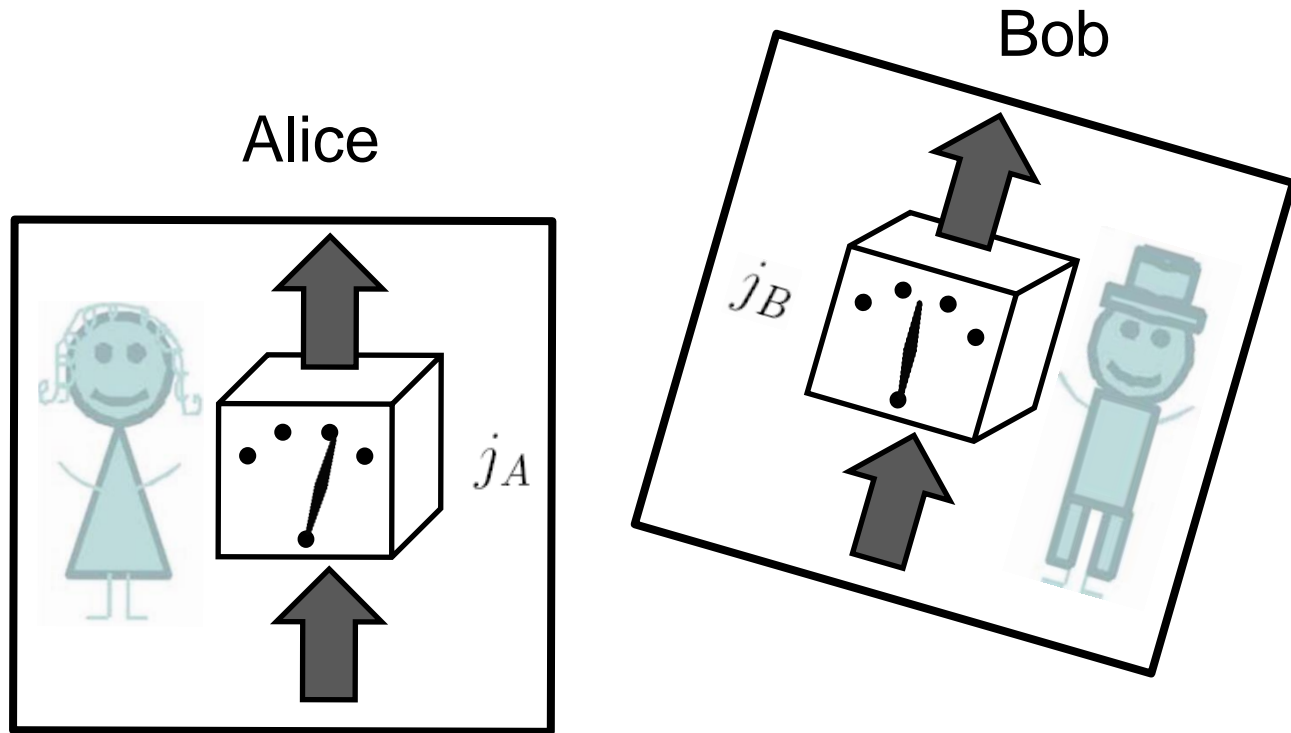
The system exits the lab.

A transformation is performed, and an outcome  $j$  is obtained.

A system enters the lab.

This is the **only** way how each party interacts with the “outside world”.

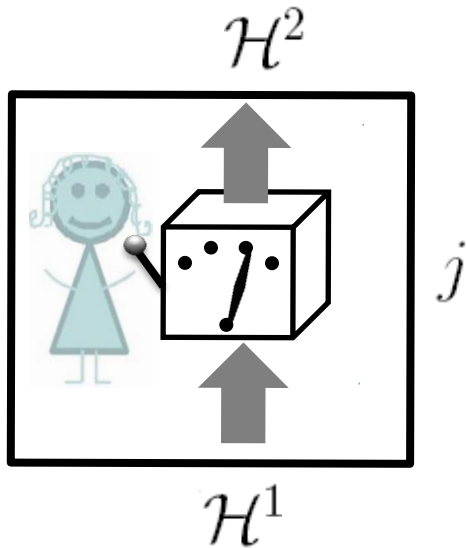
# Correlations



No prior assumption of pre-existing causal structure, in particular of the pre-existing background time.

# Main premise:

## Local descriptions agree with quantum mechanics



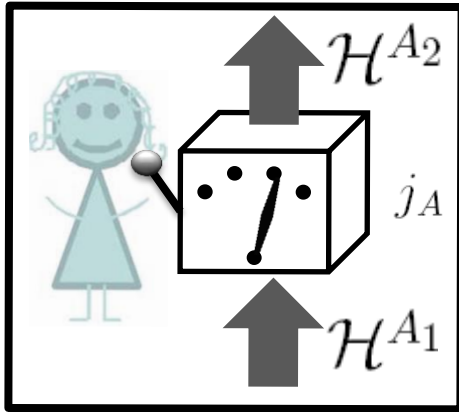
Transformations = **completely positive** (CP)  
trace-nonincreasing maps

$$\mathcal{M}_j : \mathcal{L}(\mathcal{H}^1) \rightarrow \mathcal{L}(\mathcal{H}^2)$$

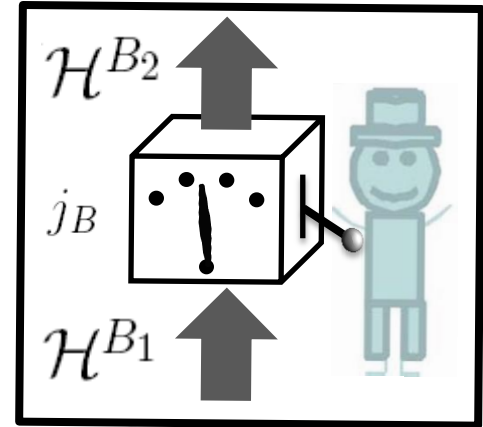
$$\mathcal{M}_{CPTP} = \sum_j \mathcal{M}_j$$

Completely positive trace preserving (CPTP) map

# Two (or more) parties



$$\mathcal{M}_{j_A}^A : \mathcal{L}(\mathcal{H}^{A_1}) \rightarrow \mathcal{L}(\mathcal{H}^{A_2})$$



$$\mathcal{M}_{j_B}^B : \mathcal{L}(\mathcal{H}^{B_1}) \rightarrow \mathcal{L}(\mathcal{H}^{B_2})$$

Probabilities are **bilinear** functions of the CP maps

$$P(\mathcal{M}^A, \mathcal{M}^B)$$

**Goal:** characterize the most general probability distributions

# Choi-Jamiołkowski isomorphism

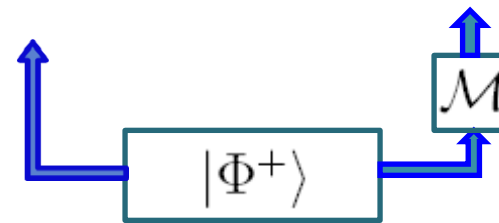
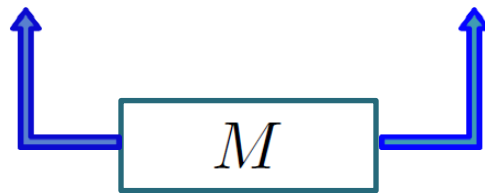
CP maps

$$\mathcal{M} : \mathcal{L}(\mathcal{H}^1) \rightarrow \mathcal{L}(\mathcal{H}^2)$$



Positive matrices

$$M \in \mathcal{L}(\mathcal{H}^1) \otimes \mathcal{L}(\mathcal{H}^2)$$



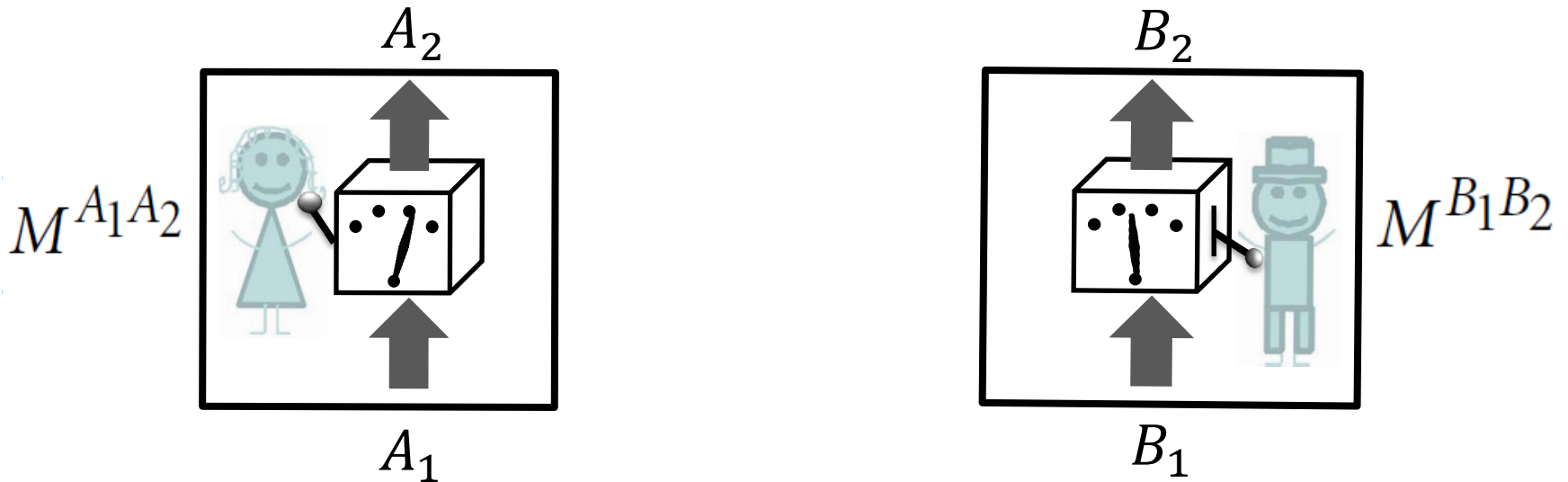
Maximally entangled state

$$|\Phi^+\rangle = \sum_i |i\rangle|i\rangle$$

$$|i\rangle \in \mathcal{H}^1$$

$$M^{12} := [\mathcal{I} \otimes \mathcal{M} (|\Phi^+\rangle\langle\Phi^+|)]^T$$

# Two (or more) parties



Probabilities are **bilinear** functions of the CP maps

$$P(\mathcal{M}^A, \mathcal{M}^B) = \text{Tr} \left[ W^{A_1 A_2 B_1 B_2} \left( M^{A_1 A_2} \otimes M^{B_1 B_2} \right) \right]$$

„Process matrix“

Choi-Jamilkowski representation of CP maps

**Goal:** characterize the most general  $W$



# Bipartite probabilities

$$P(\mathcal{M}^A, \mathcal{M}^B) = \text{Tr} \left[ W^{A_1 A_2 B_1 B_2} \left( M^{A_1 A_2} \otimes M^{B_1 B_2} \right) \right]$$

## 1. Nonnegative probabilities:

(ancillary entangled states do not fix causal order)

$$W^{A_1 A_2 B_1 B_2} \geq 0$$

## 2. Probability is 1 for all CPTP maps.

$$\text{Tr} \left[ W^{A_1 A_2 B_1 B_2} \left( M_{\text{CPTP}}^{A_1 A_2} \otimes M_{\text{CPTP}}^{B_1 B_2} \right) \right] = 1,$$

# Terms appearing in process matrix

$$W^{A_1 A_2 B_1 B_2} = \sum_{\mu_1, \dots, \mu_4} a_{\mu_1 \dots \mu_4} \sigma_{\mu_1}^{A_1} \otimes \dots \otimes \sigma_{\mu_4}^{B_2}$$

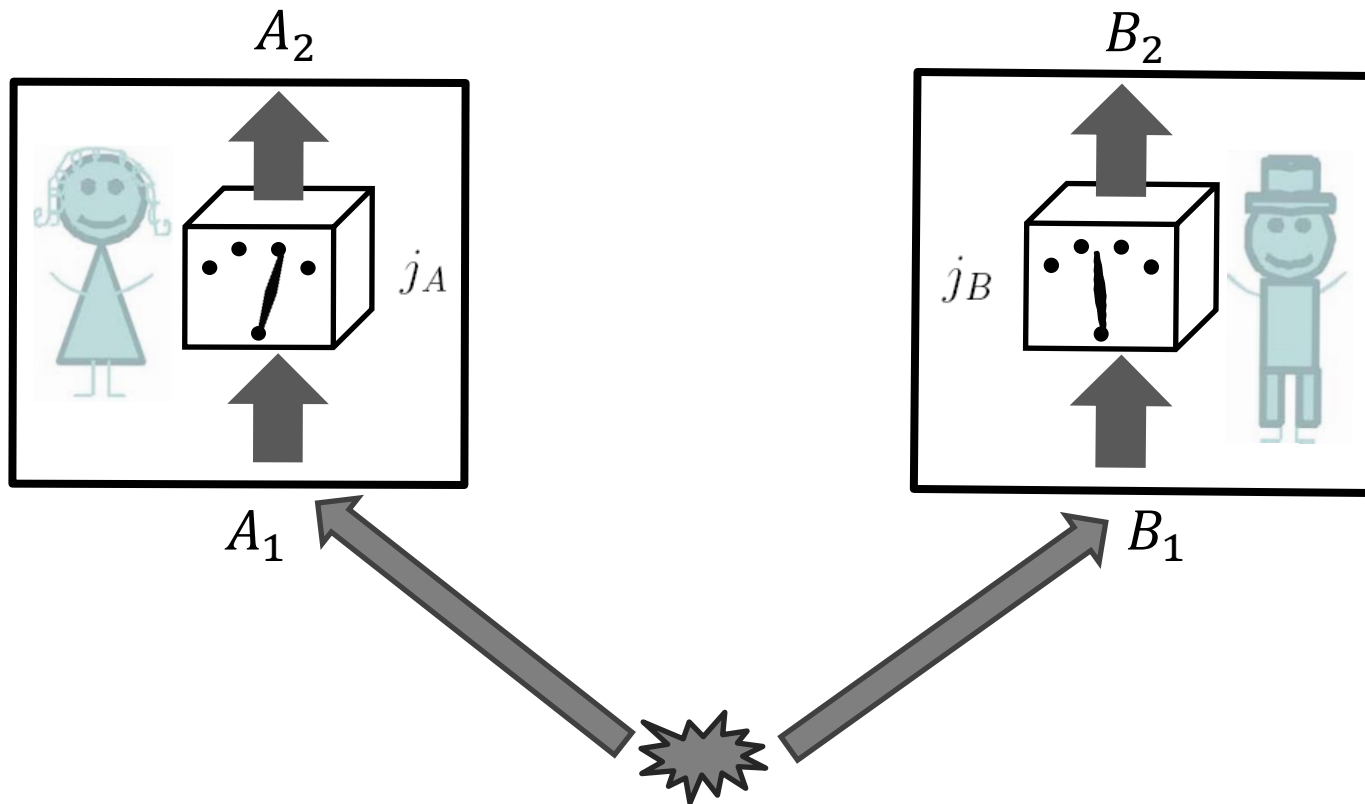
$$\sigma_i^{A_1} \otimes \mathbb{1}^{rest} \quad \text{type } A_1$$

$$\sigma_i^{A_1} \otimes \sigma_j^{A_2} \otimes \mathbb{1}^{rest} \quad \text{type } A_1 A_2$$

...

	$A_2 B_1$	$A_1 A_2 B_1$
	$A_1 B_2$	$A_1 B_1 B_2$
States	Channels	Channels with memory

# Example: Bipartite state

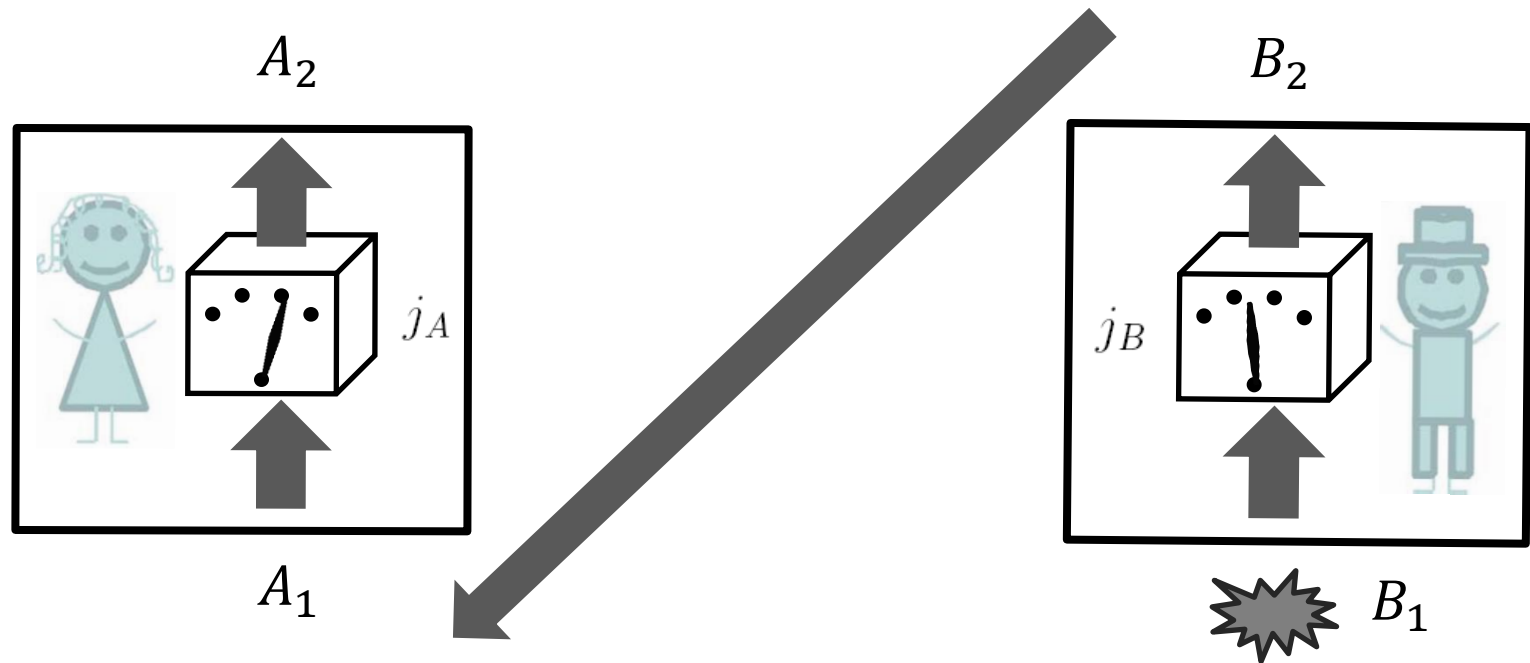


Sharing a joint state:  
(no signalling)

$$W^{A_1 A_2 B_1 B_2} = \rho^{A_1 B_1} \otimes \mathbb{1}^{A_2 B_2}$$

$$P(\mathcal{M}^A, \mathcal{M}^B) = \text{Tr} [\mathcal{M}^A \otimes \mathcal{M}^B (\rho^{A_1 B_1})]$$

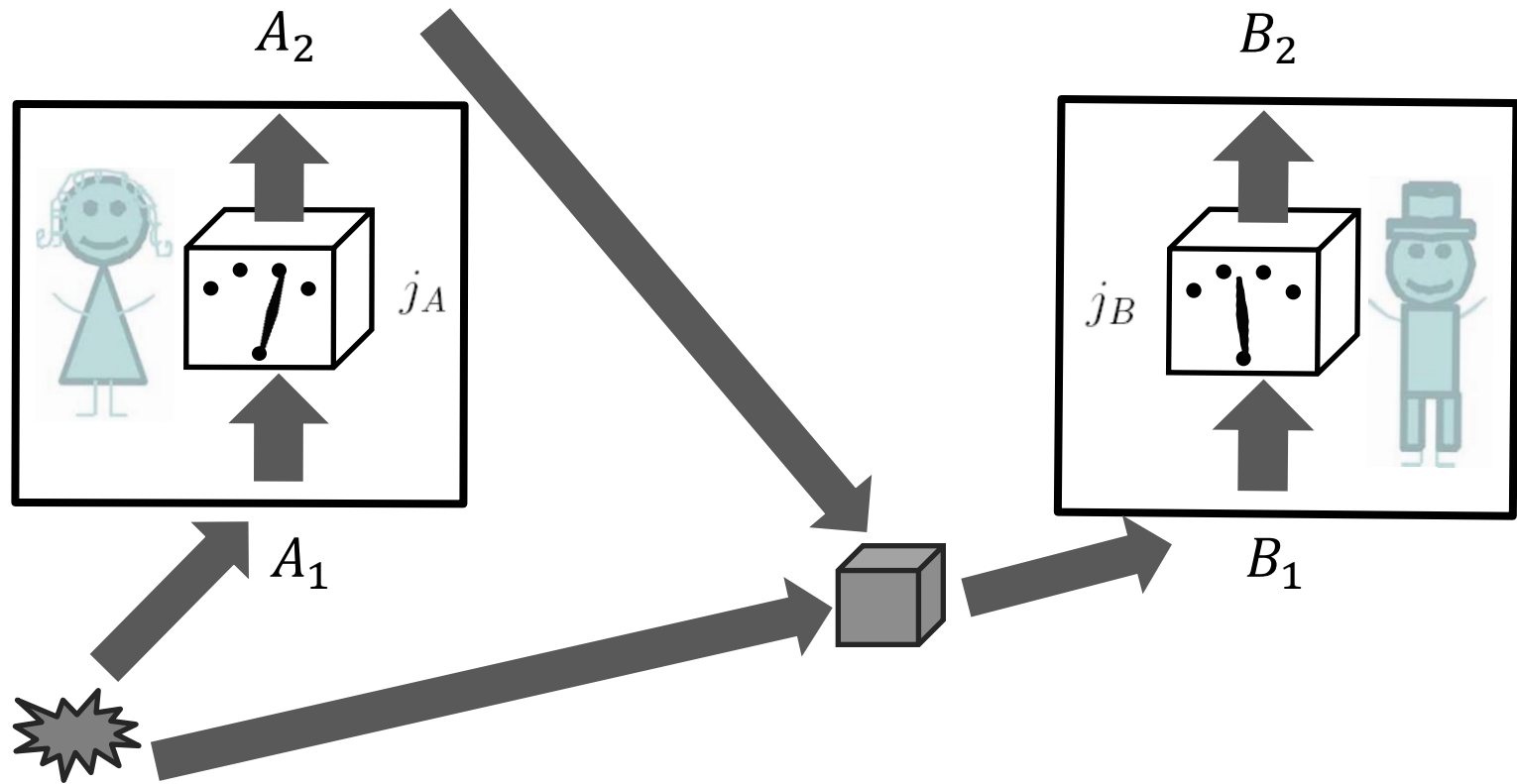
# Example: Channel $B \rightarrow A$



Sending a state from B to A:  $W^{A_1 A_2 B_1 B_2} = \mathbb{1}^{A_2} \otimes (C^{A_1 B_2})^T \otimes \rho^{B_1}$   
 (possibility of signalling)  $P(\mathcal{M}^A, \mathcal{M}^B) = \text{Tr} [\mathcal{M}^A \circ \mathcal{C} \circ \mathcal{M}^B (\rho^{B_1})]$

# Channel with memory

The most general possibility in a causal structure



Possibility of signalling:

$$W^{A_1 A_2 B_1 B_2} = W^{A_1 A_2 B_1} \otimes \mathbb{1}^{B_2}$$

# Causal order between parties

$W^{B \not\prec A}$  Bob cannot signal to Alice

$W^{A \not\prec B}$  Alice cannot signal to Bob

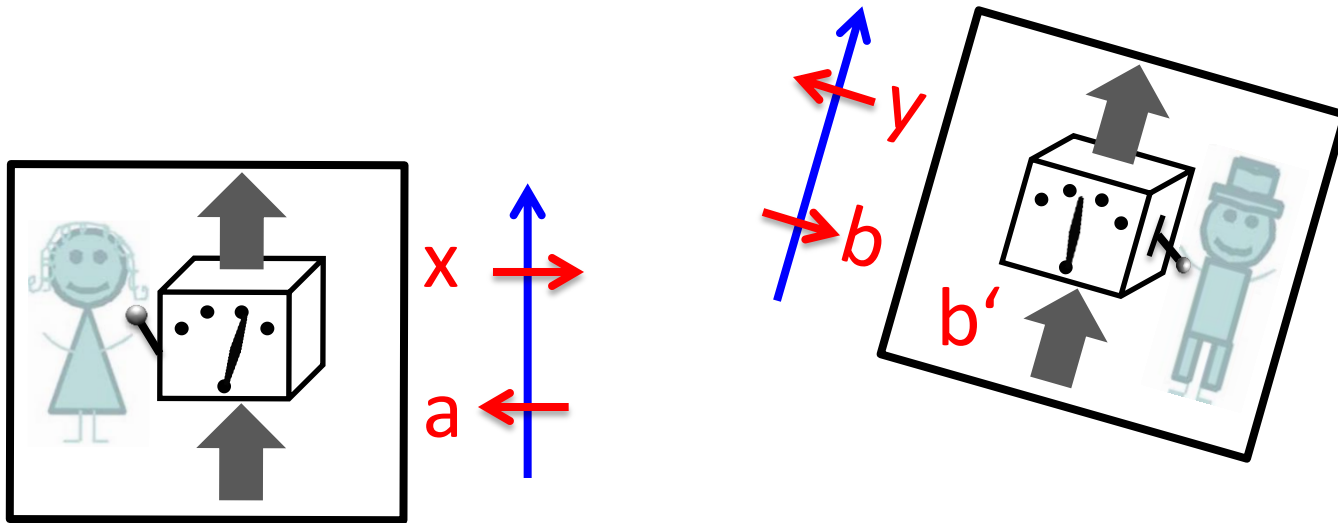
Probabilistic mixtures of ordered processes:

$$W^{A_1 A_2 B_1 B_2} = q W^{B \not\prec A} + (1 - q) W^{A \not\prec B}$$

Are all  $W$  of that form?

**No!**

# Causal game: Guess partner's input



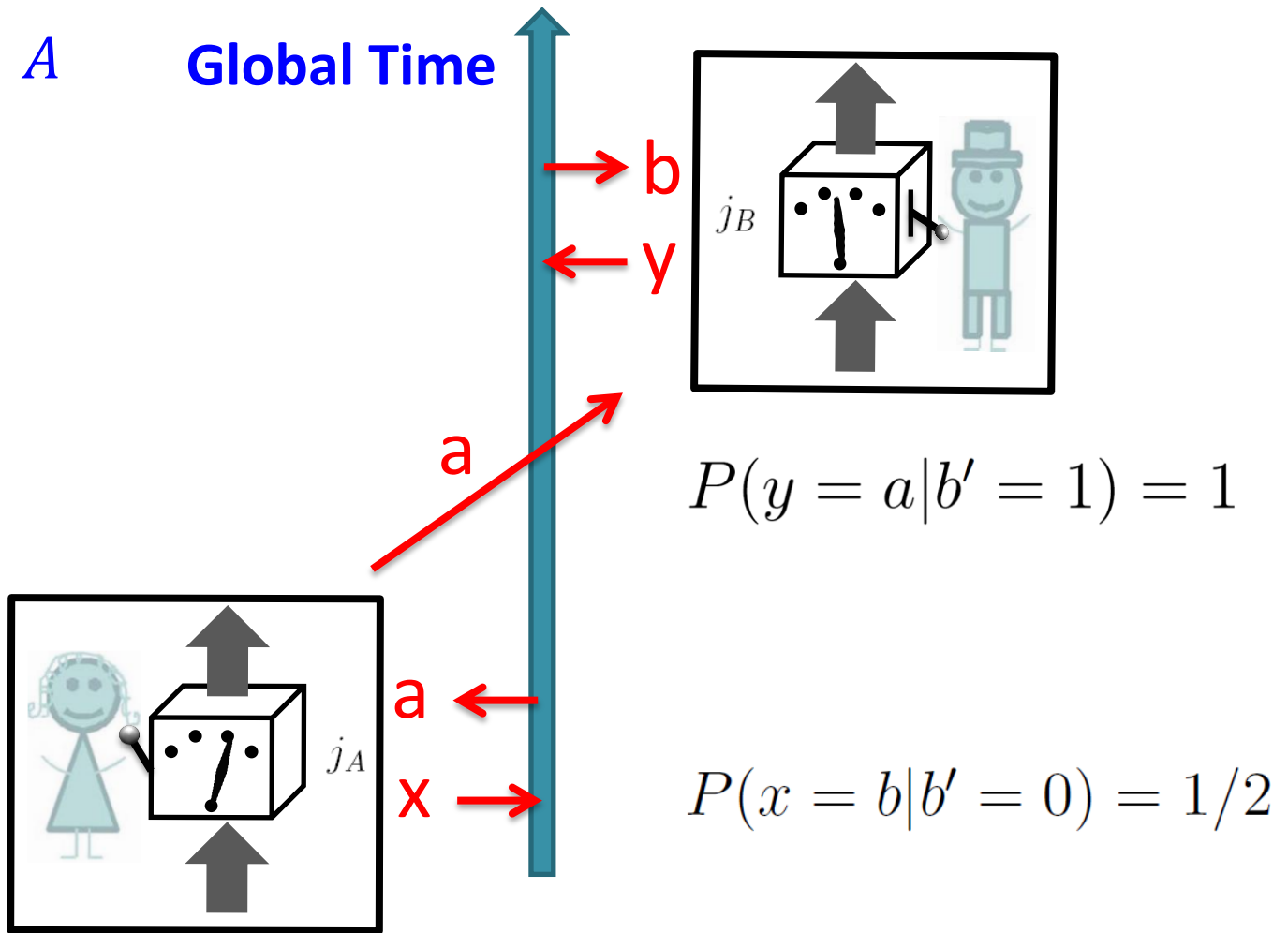
- Alice is given bit **a** and Bob bit **b**.
- Alice produces **x** and Bob **y**, which are their best guesses for the value of the bit given to the other.
- Bob is given an additional bit **b'** that tells him whether he should guess her bit (**b'=1**) or she should guess his bit (**b'=0**).
- The goal is to maximize the probability for correct guess:

$$p_{succ} := \frac{1}{2} [P(x = b|b' = 0) + P(y = a|b' = 1)]$$

# Causally ordered situation

Case:  $B \not\prec A$

Global Time



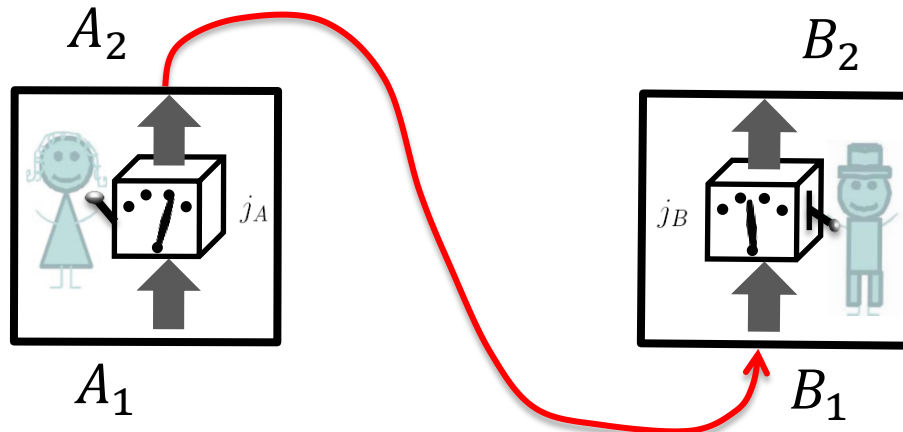
$$P_{succ} := \frac{1}{2} [P(x = b | b' = 0) + P(y = a | b' = 1)] \leq \frac{3}{4}$$



# Indefinite Causal Structures

$$W^{A_1 A_2 B_1 B_2} = \frac{1}{4} \left[ \mathbb{1}^{A_1 A_2 B_1 B_2} + \frac{1}{\sqrt{2}} \left( \sigma_z^{A_2} \sigma_z^{B_1} + \sigma_z^{A_2} \sigma_x^{B_2} \right) \right]$$

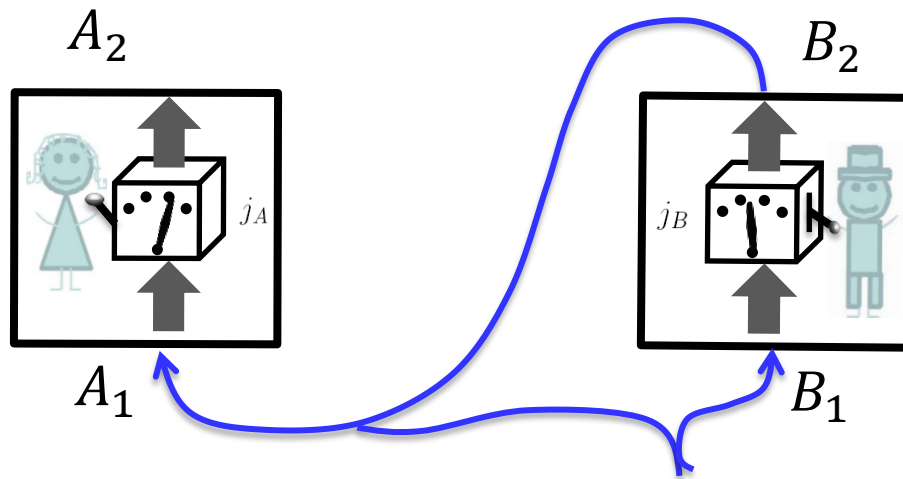
← two-level systems



**$b'=1$ :** Bob measures  $\sigma_z^{B_1}$

# Indefinite Causal Structures

$$W^{A_1 A_2 B_1 B_2} = \frac{1}{4} \left[ \mathbb{1}^{A_1 A_2 B_1 B_2} + \frac{1}{\sqrt{2}} \left( \cancel{\sigma_z^{A_2} \sigma_z^{B_1}} + \sigma_z^{A_1} \sigma_x^{B_1} \sigma_z^{B_2} \right) \right]$$

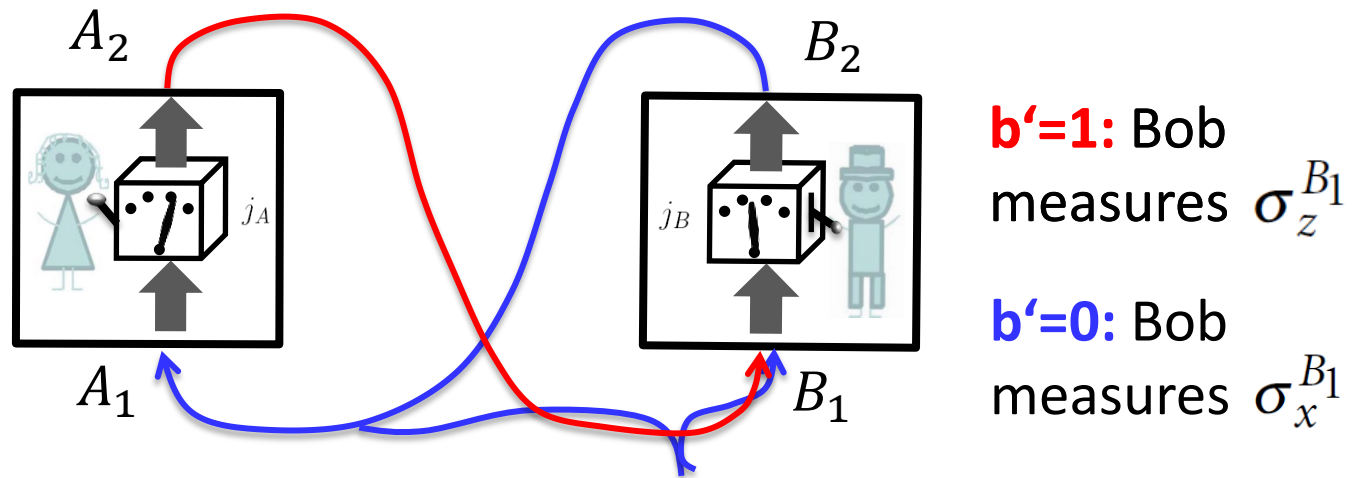


**$b'=1$ :** Bob measures  $\sigma_z^{B_1}$

**$b'=0$ :** Bob measures  $\sigma_x^{B_1}$

# Indefinite Causal Structures

$$W^{A_1 A_2 B_1 B_2} = \frac{1}{4} \left[ \mathbb{1}^{A_1 A_2 B_1 B_2} + \frac{1}{\sqrt{2}} \left( \sigma_z^{A_2} \sigma_z^{B_1} + \sigma_z^{A_1} \sigma_x^{B_1} \sigma_z^{B_2} \right) \right]$$



Depending on his choice Bob can end up **after** or **before** Alice with probability  $\sqrt{2}/2$

**The probability of success is:**  $p_{succ} = \frac{2+\sqrt{2}}{4} > \frac{3}{4}$

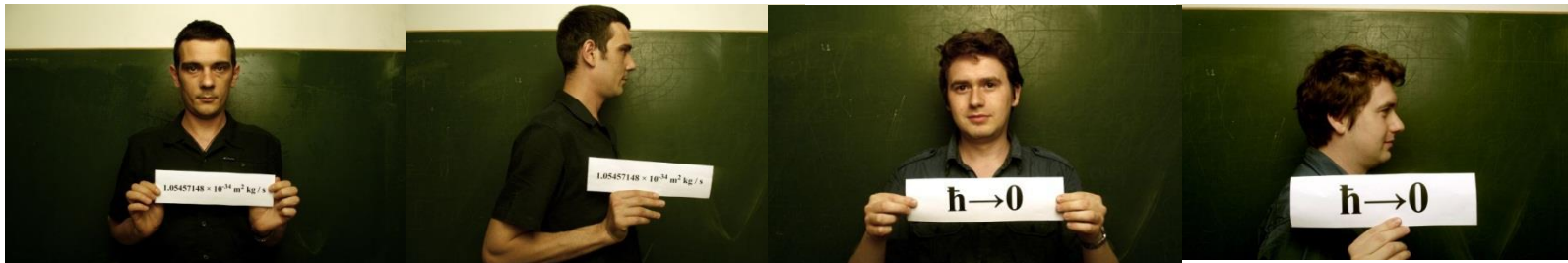
# Summary

- Unified framework for signalling (time-like) and no-signalling (space-like) quantum correlations.
- Situations where the causal order between laboratory operations is not definite → global causal order need not be a necessary element of quantum theory.
- In the classical limit, global causal order emerges.

# Outlook

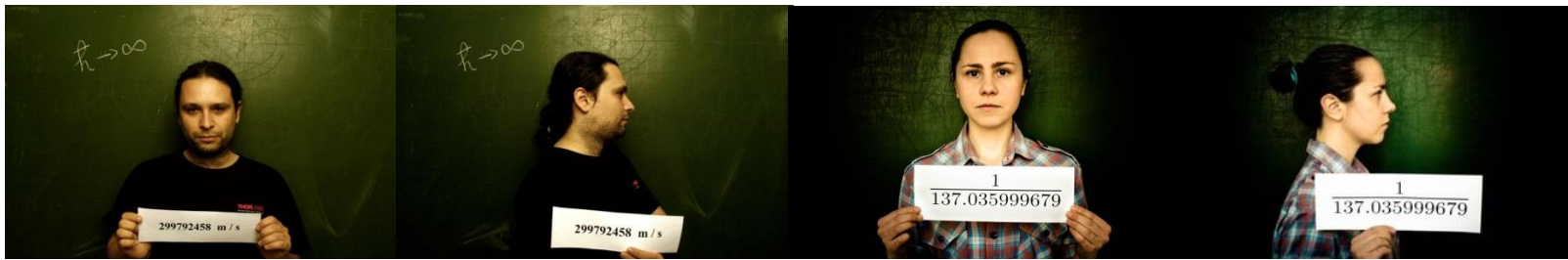
- Can we realize non-causal processes in the lab?
- A generalization of concept of space-time?
- Principles that select the generally signalling correlations allowed by QM?
- Is  $\frac{2 + \sqrt{2}}{4}$  a “Tsirelson bound for non-causal correlations”?
- A new resource for quantum information processing?

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C.B.



CoQuS

Complex  
Quantum  
Systems

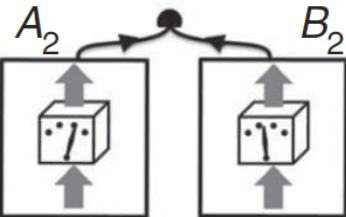
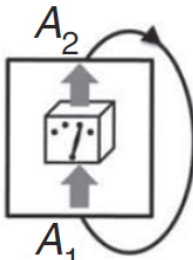
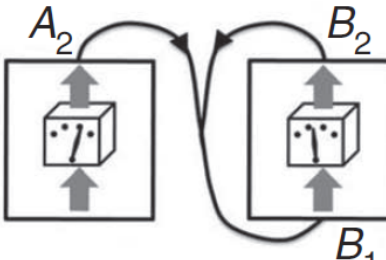
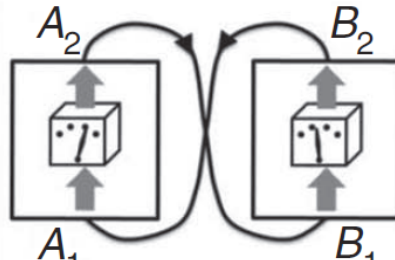


Thank you for your attention

# Terms **not** appearing in process matrix

$$W^{A_1 A_2 B_1 B_2} = \sum_{\mu_1, \dots, \mu_4} a_{\mu_1 \dots \mu_4} \sigma_{\mu_1}^{A_1} \otimes \dots \otimes \sigma_{\mu_4}^{B_2}$$

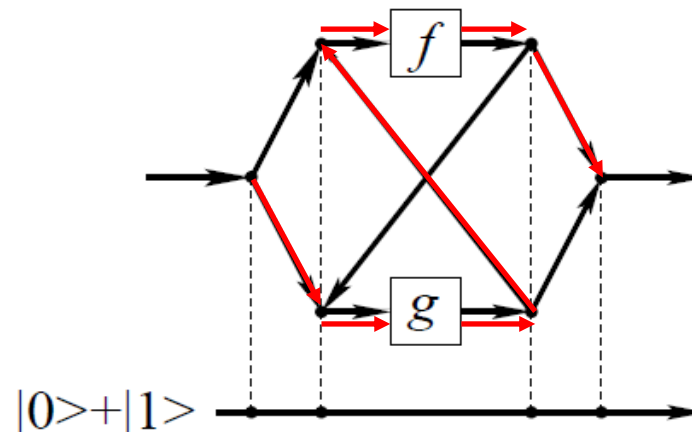
$$\begin{aligned} & \sigma_i^{A_1} \otimes \mathbb{1}^{rest} && \text{type } A_1 \\ & \sigma_i^{A_1} \otimes \sigma_j^{A_2} \otimes \mathbb{1}^{rest} && \text{type } A_1 A_2 \\ & \dots && \end{aligned}$$

$A_2, B_2, A_2 B_2$	$A_1 A_2, B_1 B_2$	$A_1 A_2 B_2, A_2 B_1 B_2$	$A_1 A_2 B_1 B_2$
Postselection	Local loops	Channels with local loops	Global loops
			

# Where to look for non-causal processes?

1. **Within standard quantum mechanics?** It may be possible to create causally non-separable situations making clever use of superposition and entanglement.

Pavia group,  
Chriribella et. al





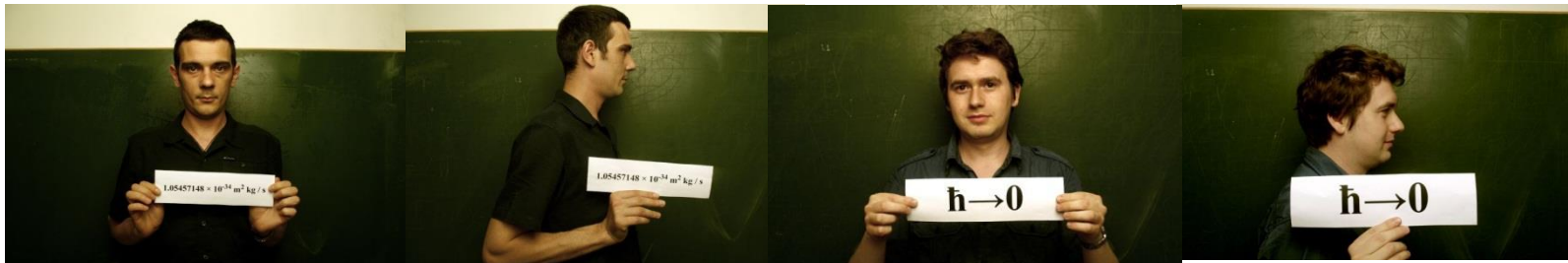
# Conclusions

- [Not shown]: In the classical limit all correlations are causally ordered
- Unified framework for both signalling (“time-like”) and non-signalling (“space-like”) quantum correlations
- Situations where a causal ordering between laboratory operations is not definite → Suggests that causal ordering might not be a necessary element of quantum theory

# Outlook

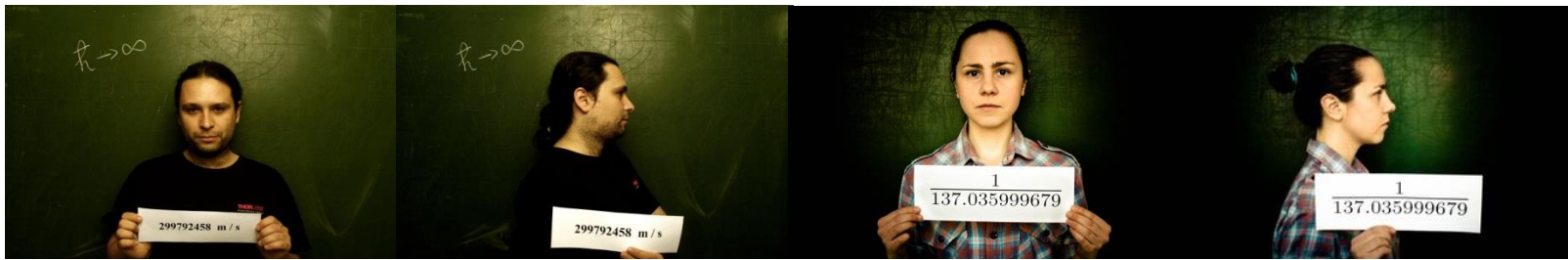
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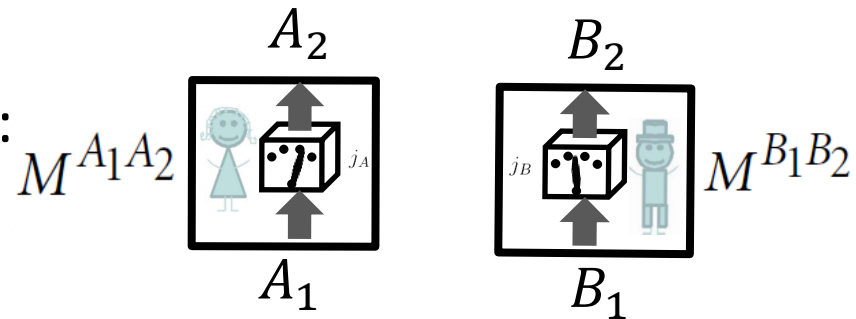
FQXi John Templeton Foundation  
FOUNDATIONAL QUESTIONS INSTITUTE

Thank you for your attention

quantumfoundations.weebly.com

# Bipartite probabilities

**Bilinear** functions of the CP maps:



Representation

$$P(\mathcal{M}^A, \mathcal{M}^B) = \text{Tr} \left[ W^{A_1 A_2 B_1 B_2} \left( M^{A_1 A_2} \otimes M^{B_1 B_2} \right) \right]$$

„**Process matrix**“

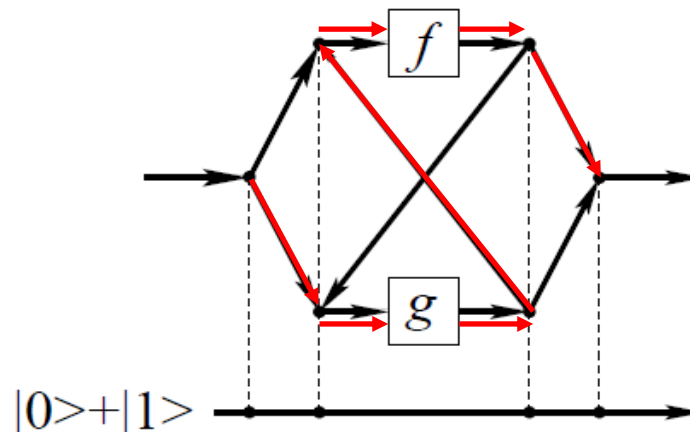
Choi-Jamilkowski representation of CP maps

**Goal:** characterize the most general  $W$

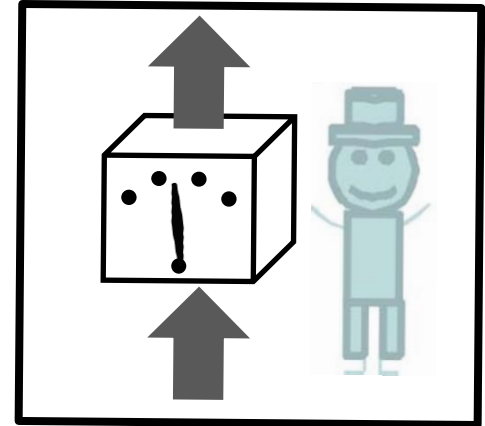
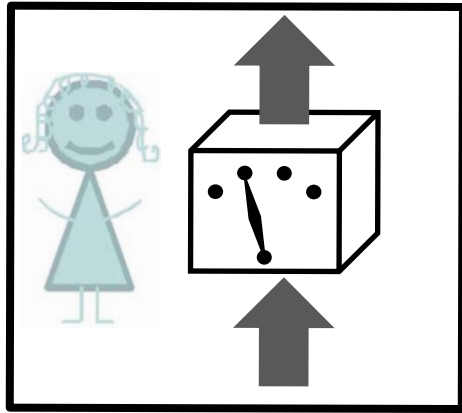
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1. **Within standard quantum mechanics?** It may be possible to create causally non-separable situations making clever use of superposition and entanglement.

Pavia group,  
Chriribella et. al



# Causal order from correlations?



“In summer with a large amount of ice-cream consumption there are lot of sun-burn cases.”



“Ice consumption causes sun-burn.”

# Motivation

- Can one formulate physical theories without the assumption of background space-time or causal structure?  
*Using tools of quantum information to address problems that traditionally have been considered within quantum gravity*
- Quantum correlations are the crucial resource for performing computational tasks that are impossible classically.  
*“Superpositions of quantum circuits”*

L. Hardy, arXiv:gr-qc/0509120, G. Chiribella, G. M. D'Ariano, P. Perinotti, B. Valiron, arXiv:0912.0195v3, M. S. Leifer, R. W. Spekkens, arXiv:1107.5849

# Indefinite Causal Structures

Alice always measures in the z basis and encodes the bit in the z basis

Alice's CP map:  $|z_x\rangle\langle z_x|^{A_1} \otimes |z_a\rangle\langle z_a|^{A_2} \quad x, a = \pm 1$

If Bob wants to receive ( $\mathbf{b}'=1$ ), he measures in the z basis



$$W^{A_1 A_2 B_1 B_2} = \frac{1}{4} \left[ \mathbb{1} + \frac{1}{\sqrt{2}} \left( \sigma_z^{A_2} \sigma_z^{B_1} + \cancel{\sigma_z^{A_1} \sigma_x^{B_1} \sigma_z^{B_2}} \right) \right]$$

$\langle z_{\pm} | \sigma_x | z_{\pm} \rangle = 0$



Channel from Alice to Bob

Not seen by Bob



Bob receives the state

$$\widetilde{W}^{B_1 B_2} = \frac{1}{2} \left( \mathbb{1} + a \frac{1}{\sqrt{2}} \sigma_z^{B_1} \right)$$

He can read Alice's bit with probability

$$P(y = a | b' = 1) = \frac{2 + \sqrt{2}}{4}$$



# Indefinite Causal Structures

If Bob wants to send ( $\mathbf{b}' = \mathbf{0}$ ), he measures in the  $x$  basis and encodes in the  $z$  basis conditioned on his outcome

Bob's CP map:  $|x_y\rangle\langle x_y|^{B_1} \otimes |z_{by}\rangle\langle z_{by}|^{B_2} \quad y, b = \pm 1$



$$W^{A_1 A_2 B_1 B_2} = \frac{1}{4} \left[ \mathbb{1} + \frac{1}{\sqrt{2}} \left( \sigma_z^{A_2} \sigma_z^{B_1} - \sigma_z^{A_1} \sigma_x^{B_1} \sigma_z^{B_2} \right) \right]$$

$$\langle x_{\pm} | \sigma_z | x_{\pm} \rangle = 0$$

Not seen by Bob

Channel from Bob to Alice, correlated with Bob's outcome



Alice receives the state

$$\widetilde{W}^{A_1 A_2} = \frac{1}{2} \left( \mathbb{1} + b \frac{1}{\sqrt{2}} \sigma_z^{A_1} \right)$$

She can read Bob's bit with probability

$$P(x = b | b' = 0) = \frac{2 + \sqrt{2}}{4}$$

# Indefinite Causal Structures

$$W^{A_1 A_2 B_1 B_2} = \frac{1}{4} \left[ \mathbb{1}^{A_1 A_2 B_1 B_2} + \frac{1}{\sqrt{2}} \left( \sigma_z^{A_2} \sigma_z^{B_1} + \sigma_z^{A_1} \sigma_x^{B_1} \sigma_z^{B_2} \right) \right]$$

There are quantum processes that **cannot** be understood as probabilistic mixtures of causally ordered situations!



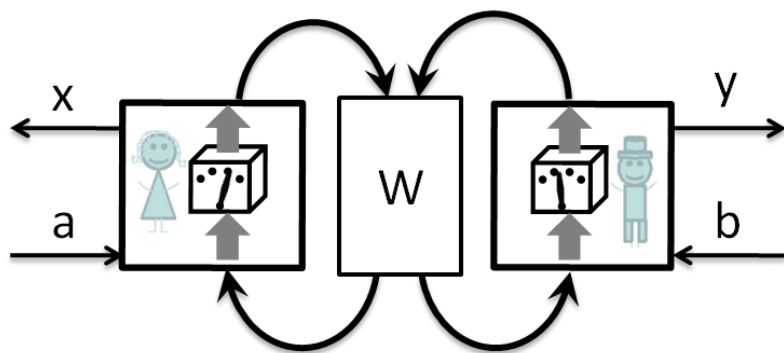
Artistic view on 'Superpositions of causal orders' inspired by M. C. Escher's *Ascending and Descending*, 1960

Jonas Schmöle, University of Vienna

# Where to look for non-causal processes?

1. **Within standard quantum mechanics?** It may be possible to create causally non-separable situations similarly to mixtures, making clever use of superposition and entanglement.

2. **Closed time-like curves?** Every  $W$  can be seen as a noisy channel back in time.



Grandfather paradox is avoided.

This CTC-like structure is linear, unlike Deutsch's or Bennet's CTC models.

# Where to look for non-causal processes?

3. “Superposition of causal orders”? Every  $W$  contains terms that correspond to situations with definite causal order, yet is not a classical mixture of those.

**[Not shown]:** In the classical limit all processes are causally separable, i.e., global causal order arises!

**Space-time may emerge in a quantum-to-classical transition.**